# A PULSED NMR STUDY OF MOLECULAR MOTION IN SOLID 5,5-DIMETHYL-5,6,11,12-TETRAHYDRO-5H-DIBENZO[b,f]SILOCIN 

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## Summary

A pulsed NMR study of relaxation times, $T_{1}$ and $T_{1 D}$, and second moments for solid 5,5-dimethyl-5,6,11,12-tetrahydro-5H-dibenzo $[b, f]$ silocin is reported. The compound was studied over the temperature range -175 to $50^{\circ} \mathrm{C}$. Evidence was obtained for three motions characerized by activation energies, 2.70, 1.2, and $\geqslant 4 \mathrm{kcal} / \mathrm{mole}$, respectively. The first motion is methyl reorientation. The second motion is not assigned but is probably a relatively small amplitude flexing of the central ring. The third motion is also unassigned and may be large amplitude flexing of the central ring, libration, or anisotropic molecular reorientation. The motions were assigned by comparison of $X$-ray crystallographic data with NMR second moment data.

## Introduction

We have previously reported the results of pulsed NMR studies of molecular motion in tricyclic compounds with central eight-membered rings in the solid phase. The compounds studied were 6,12,12-trimethyl-5,6-dihydro-7H,12Hdibenzo [c,f][1,5]silazocine [1] and $N$-methyl-5,6-dihydro-7H,12H-dibenz [c,f]azocine [2]. In the solid phase, the silazocine adopts a flexible, distorted twist hoat (TB) conformation and the azocine adopts the rigid, boat-chair (BC) conformation. The motions exhibited by each compound were found to depend upon the conformation of the central 8-membered ring. Recently, a dibenzosilocin was synthesized [3] and its X-ray structure was reported [4]. This compound adopts a basket conformation in the solid phase, and examination of molecular models suggests that the central ring of this compound may exhibit
some degree of flexibility. We have performed a pulsed NMR study of this compound to determine the degree of flexibility of the central ring and to compare the results from this compound with those reported previously.

## Experimental

The silocin, 5,5 -dimethyl-5,6,11,12-tetrahydro-5H-dibenzo[b,f]silocin (I), the subject of this study, was prepared by the ring expansion reaction of 5 -chloromethyl-5-methyl-10,11-dihydrodibenzo [b,f]silepin in the presence of $\mathrm{AlCl}_{3}$ followed by quenching with methyllithium. Distillation of the crude reac-

tion product followed by recrystallization from ethanol gave pure silocin, m.p. $60-70.5^{\circ} \mathrm{C} *$. The sample was placed in a glass container for study. The proton NMIR measurements [5] ** of the Zeeman spin lattice relaxation time ( $T_{1}$ ), the dipolar relaxation time ( $T_{\text {ID }}$ ), and the second moment ( $M_{2}$ ) were made using a Polaron (Watford, England) high power pulsed NMR spectrometer operating at 60 MHz , as described previously [1]. The second moment values were obtained from analysis of the "solid echo" following a $90-\tau-90$ pulse sequence [6,7], in which the second pulse is phase shifted by $90^{\circ}$ from the first pulse. The Bloch decays were found to be Gaussian within experimental error for this polycrystalline sample.

Computer calculations were done on an IBM 370/168 computer.

## Results

Experimental values of the Zeeman spin-lattice relaxation time $\left(T_{1}\right)$ and the dipolar relaxation time ( $T_{1 \mathrm{D}}$ ) are presented in Fig. 1. $T_{1}$ exhibits a relaxation minimum centered at $10^{3} / T=7.3(\sim 138 \mathrm{~K})$ due to molecular motion (motion ). The solid line for $T_{1}$ in Fig. 1 represents values for $T_{1}$ calculated from eq. 1.

$$
\begin{equation*}
\frac{1}{T_{1}}=\frac{2}{3} \gamma^{2} M_{2 \bmod }\left[\frac{\tau_{c}}{1+\omega_{o}^{2} \tau_{c}^{2}}+\frac{4 \tau_{c}}{1+4 \omega_{o}^{2} \tau_{c}^{2}}\right] \tag{1}
\end{equation*}
$$

Where $\gamma$ is the gyromagnetic ratio for protons, $M_{2 \text { mod }}$ is the portion of the second moment $M_{2}$ which is modulated by the motion, $\omega_{0}$ is the Larmor frequency in the laboratory magnetic field, and $\tau_{c}$ is the correlation time of the motion. The line for $T_{1}$ in Fig. 1 is calculated assuming an Arrhenius expression

[^0]

Fig. 1. Observed and calculated values of spin-lattice and dipolar relaxation times for dibenzsilocin vs. reciprocal temperature.
(eq. 2), in which $E_{\mathrm{A}}=2.5 \mathrm{kcal} / \mathrm{mole}$ and $\tau_{\mathrm{o}}=2.1 \times 10^{-13} \mathrm{sec}$. It can be seen in $\tau_{\mathrm{c}}=\tau_{\mathrm{o}} \exp \left(E_{\mathrm{A}} / R T\right)$
Fig. 1 that the calculated $T_{1}$ values deviate from the experimental values on the low temperature side of the minimum. The gradient of the observed $T_{1}$ values for $10^{3} / T>8$ corresponds to $E_{\mathrm{A}} \sim 1.7 \mathrm{kcai} / \mathrm{mole}$. In addition, there is some discrepancy between observed and calculated $T_{1}$ values for $10^{3} / T<3.5$.

For $10^{3} / T>4$, values of $T_{10}$ are dominated by the same process (motion I) that controls $T_{1}$. For $10^{3} / T<4$, there is an indication that an additional process (motion II) contributes to $T_{1 \mathrm{D}}$. The solid line for $T_{1 \mathrm{D}}$ in Fig. 1 is calculated from an expression of Goldman [8], which assumes completely uncorrelated motion (eqn. 3).
$T_{1 \mathrm{D}}=T_{1} /\left(2+\omega_{\mathrm{o}}^{2} \tau_{\mathrm{c}}^{2} / 3\right)$
It can be seen in Fig. 1 that there is good agreement between calculated and observed values of $T_{1 \mathrm{D}}$ for $10^{3} / T>4$. The dominant process (motion I) controlling $T_{1}$ and $T_{10}$ in Fig. 1 is random reorientation of both methyl groups; as discussed below.

For $10^{3} / T<4$, both motion I and motion II contribute to the observed value of $T_{1 \mathrm{D}}$ according to eq. 4.

$$
\begin{equation*}
\left(\frac{1}{T_{1 \mathrm{D}}}\right)_{\mathrm{obs}}=\left(\frac{1}{T_{1 \mathrm{D}}}\right)_{\mathrm{I}}+\left(\frac{1}{T_{1 \mathrm{D}}}\right)_{\mathrm{II}} \tag{4}
\end{equation*}
$$

Values of ( $\left.T_{1 \mathrm{D}}\right)_{\text {II }}$ obtained from eq. 4 are presented as a dashed line in Fig. 1. The gradient of this line corresponds to $\left(E_{\mathrm{A}}\right)_{\mathrm{II}}=4.1 \mathrm{kcal} / \mathrm{mole}$; however, this
can be taken only as a rough lower limit of the activation energy for motion II.
Thus, values of $T_{1 D}$ for $10^{3} / T>4$ appear to be controlled by a single relaxation process, motion $I$, (methyl reorientation), whereas values of $T_{1}$ appear to be controlled principally by motion $I$ with an indication of a contribution from at least one additional process (motion III). The discrepancy between observed and calculated $T_{1}$ values for $10^{3} / T<3.5$ cannot be attributed to a contribution to $T_{1}$ from motion II, since the contribution to $T_{10}$ from motion $\Pi$ in this temperature range indicates that the correlation frequency of motion II is $<10 \mathrm{~Hz}$ for $10^{3} / T>3$, and thus the contribution to $T_{1}$ from motion II is negligible.

The behavior of $T_{1}$ over the entire range of temperature studied, as shown in Fig. 1, can be accounted for with two expressions of the form of eq. 1, corresponding to motion I and III. Values of $T_{1}$ so calculated are shown in Fig. II, and it can be seen that there is good agreement between observed and calculated $T_{1}$ values over the entire temperature range. The parameters found for motion I are: $E_{\mathrm{A}}=2.70 \mathrm{kcal} / \mathrm{mole}, \tau_{\mathrm{o}}=8.6 \times 10^{-14} \mathrm{sec}$, and $M_{2 \mathrm{mod}}=2.1 \mathrm{G}^{2}$; the parameters found for motion II are: $E_{\mathrm{A}}=1.15 \mathrm{kcal} / \mathrm{mole}, \tau_{\mathrm{o}}=9.2 \times 10^{-12} \mathrm{sec}$, and $M_{2_{\text {rad }}}=0.25 G^{2}$. Thus motion III modulates a much smaller portion of the total dipolar interaction than does motion $I$.

It should be noted that the qualitative shape of the methyl reorientation. minimum in Fig. 1 and 2 may also be accounted for with the assumption of a temperature dependent distribution of correlation frequencies, such as is postulated for polymers [9]. However, such a circumstance is most probable for an amorphous material and is less likely in a polycrystalline sample. It should also be noted that the curves in Fig. 2 do not represent a unique fit of the experimental data. The curve for motion III is somewhat arbitrary; however, in attempting to fit the data, we conclude (a) the $T_{1}$ minimum for motion III must be near its position in Fig. 2 and (b) motion III must have an activation energy about $1 \mathrm{kcal} / \mathrm{mole}$. The two curves in Fig. 2 represent the simplest way


Fig. 2. Observed and calcuiated values of spin-latice relaxation times for dibenzsilocin vs reciprocal temperature.
to account for both low temperature and high temperature discrepancies between observed and calculated $T_{1}$ values in Fig. 1.

## Discussion

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The second moment is anticipated [1,2] to be essentially constant over the entire temperature range show in Fig. 1. Experimentally we observe $M_{2}=12.0$ $\mathrm{G}^{2}$ at $22^{\circ} \mathrm{C}$ and $12.6 \mathrm{G}^{2}$ at $-173^{\circ} \mathrm{C}$. This value of $M_{2}$ corresponds to $M_{2(\text { high })}$ with respect to motions I and III, and it corresponds to $M_{2(10 w)}$ with respect to motion II. $M_{2(\text { high }}$ is the limiting value of $M_{2}$ when the motion is sufficiently rapid to cause the NMR line to be motionally narrowed, and $M_{2 \text { (low) }}$ is the value of $M_{2}$ when the motion is sufficiently slow so as to have no effect on the NMR line. It is most likely that the principal motion (I) controlling $T_{1}$ in Fig. 1 and 2 is methyl group reorientation, and to confirm this assignment, we have calculated $M_{2}$ values from X-ray crystallographic data [4,10].

The "rigid lattice" value of $M_{2}$ is that value observed only if all motion is sufficiently slow so that the NMR line is not motionally narrowed, and this value can be calculated from the equation of Van Vleck [11] (eq. 5), where
$M_{2}=\frac{715.9}{n} \sum_{j>k} r_{j k}{ }^{-6}$
$M_{2}$ is in Gauss ${ }^{2}, n$ is the number of protons in the sample, and $r_{j k}$ is the distance between protons $j$ and $k$ in $\AA$, as obtained from crystallographic data. The sum is taken over all proton pairs in the sample, but for convenience, $M_{2}$ is separated into intramolecular and intermolecular parts. The result for $M_{2 \text { intra }}$ is $24.2 \mathrm{G}^{2}$. The value of the intermolecular contribution was calculated using a computer. The compound crystallizes in the space group Pbca, with 8 molecules/unit cell. The calculation included the unit cell and all 26 adjacent unit cells. The resulting value of $M_{2 \text { inter }}$ is $1.9 \mathrm{G}^{\mathbf{2}}$, and this yields a total value for the rigid lattice $M_{2(\text { r.1. })}=26.1 \mathrm{G}^{2}$.

The observed $M_{2}$ is less than the rigid lattice $M_{2}$ due to molecular motion, and its value can be calculated for an assumed motion using the expressions of Andrew and Eades [12] as described previously [1,2]. The results were: $M_{2 \text { intra }}$ $=13.4 \mathrm{G}^{2}, M_{2 \text { inter }}=1.4 \mathrm{G}^{2}$, and a total $M_{2(\mathrm{Me} \mathrm{Rot})}=14.8 \pm 0.6 \mathrm{G}^{2}$. We assume that the uncertainty in this value is due primarily to random error in the crystal lographic data. This may be compared with the observed $M_{2}$ value $12.0-12.6 \mathrm{G}^{2}$, with an experimental uncertainty of $0.5 \mathrm{G}^{2}$. It can be seen that the observed $M_{2}$ appears to be beyond experimental uncertainty, about $1 \mathrm{G}^{2}$ smaller than that calculated assuming only methyl reorientation. Thus, the largest reduction in the second moment ( $\sim 11.3 G^{2}$ ) and the principal relaxation minimum in Fig. 2 may be attributed to methyl group reorientation (motion I). An additional smaller reduction in the second moment ( $\sim 1 G^{2}$ ) and the weaker relaxation minimum in Fig. 2 may be attributed to motion III.

It is difficult to assign motion III since it modulates only a very small portion of the dipolar interaction; however, it is likely that this motion involves the central ring. Molecular models indicate that any flexing of the central ring would probably be relatively low in amplitude due to steric interaction between methyl and ring protons. The X-ray study [4] indicates that the thermal parameters for

TABLEI
OBSERVED MOTIONS FOR DIBENZSILOCIN

| Motion | Tentative assignment | $\begin{aligned} & E_{\mathrm{A}} \\ & \text { (kcal/mole) } \end{aligned}$ | $\begin{aligned} & \tau_{\mathbf{o}} \\ & (\mathrm{sec}) \end{aligned}$ |
| :---: | :---: | :---: | :---: |
| I | $\mathrm{CH}_{3}$ reorient. | $2.70{ }^{\text {a }}$ | $8.6 \times 10^{-14}$ |
| II | - | $\geqslant 4$ | - |
| III | (Ring flex.) | $1.2{ }^{6}$ | $9 \times 10^{-12}$ |

${ }^{a}$ Estimated experimental error $\pm 0.2 \mathrm{kcal} / \mathrm{mole}{ }^{b}$ Estimated experimental error $\pm 0.4 \mathrm{kcal} / \mathrm{mole}$.
the methyl groups are relatively large. Thus, motion III could involve a low amplitude flexing of the central ring or a more localized flexing in the vicinity of the silicon atom. A large amplitude flexing, as was observed for the silazocine central ring [1], may definitely be ruled out in this case.

It is not possible to assign motion II since we have only a rough estimate of the activation energy and no second moment data. Motion II may be a molecular libration [13], flexing of the central ring (larger amplitude than motion III), or molecular reorientation (as discussed in ref. 1 and 2).

The activation parameters for the dibenzosilocin are summarized in Table 1. The $E_{\mathrm{A}}$ value for methyl orientation is essentially the same as previously reported barriers for rotation about $\mathrm{Si}-\mathrm{CH}_{3}$ bonds $[1,14], 2.2-2.4 \mathrm{kcal} / \mathrm{mole}$.

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[^0]:    * Details of the synthetic procedures will be published at a later date.
    * Definitions of the NMR parameters ind brief discussions of theory are also given in ref. 1.

